# Lecture 3: Wasserstein Space

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The material of today's lecture is adapted from Q. Mérigot's lecture notes and [3, 4].

## 1 Reminders

Let X, Y be compact metric spaces,  $c \in \mathcal{C}(X \times Y)$  the cost function and  $(\mu, \nu) \in \mathcal{P}(X) \times \mathcal{P}(Y)$  the marginals. In previous lectures, we have seen that the optimal transport problem can be formulated as an optimization over the space of transport plans  $\Pi(\mu, \nu)$  — the primal or Kantorovich problem — and as an optimization over potential functions  $\{(\varphi, \psi) \in \mathcal{C}(X) \times \mathcal{C}(Y) \mid \varphi \oplus \psi \leq c\}$  — the dual problem. We recall the following results:

- minimizer/maximizers exist for both problems and, for the dual, can be chosen as  $(\varphi, \varphi^c)$  with  $\varphi$  c-concave.
- at optimality, it holds  $\varphi(x) + \psi(y) = c(x, y)$  for  $\gamma$ -almost every (x, y)
- we have the following special cases:
  - for  $X = Y \subset \mathbb{R}$  and c(x, y) = h(y x) with h strictly convex, the optimal transport plan is the (unique) monotone plan, which can be characterized with the quantile functions of  $\mu$  and  $\nu$ .
  - for X = Y and c(x, y) = dist(x, y), we have the Kantorovich-Rubinstein formula

$$\mathcal{T}_c(\mu,\nu) = \sup_{\varphi \text{ 1-Lip}} \int \varphi d(\mu-\nu).$$

- for  $X = Y \subset \mathbb{R}^d$  and  $c(x, y) = \frac{1}{2}|y - x|^2$ , and when  $\mu$  is absolutely continuous, there exists a unique optimal transport plan. It is of the form  $\gamma = (\mathrm{id}, \nabla \tilde{\varphi})_{\#} \mu$ for some  $\tilde{\varphi} \in \mathcal{C}(\mathbb{R}^d)$  convex.

## 2 Wasserstein space

#### 2.1 Definition and elementary properties

**Definition 2.1** (Wasserstein space). Let (X, dist) be a compact metric space. For  $p \ge 1$ , we denote by  $\mathcal{P}_p(X)$  the set of probability measures on X endowed with the *p*-Wasserstein distance, defined as

$$W_p(\mu,\nu) := \left(\min_{\gamma \in \Pi(\mu,\nu)} \int \operatorname{dist}(x,y)^p \mathrm{d}\gamma(x,y)\right)^{1/p} = \mathcal{T}_{\operatorname{dist}^p}(\mu,\nu)^{\frac{1}{p}}.$$

This distance is a natural way to build a distance on  $\mathcal{P}(X)$  from a distance on X. in particular, the map  $\delta : X \to \mathcal{P}_p(X)$  mapping a point  $x \in X$  to the Dirac mass  $\delta_x$  is an isometry.

#### **Proposition 2.2.** $W_p$ satisfies the axioms of a distance on $\mathcal{P}_p(x)$ .

*Proof.* The symmetry of the Wasserstein distance is obvious. Moreover,  $W_p(\mu, \nu) = 0$  implies that there exists  $\gamma \in \Pi(\mu, \nu)$  such that  $\int \operatorname{dist}^p d\gamma = 0$ . This implies that  $\gamma$  is concentrated on the diagonal, so that  $\gamma = (\operatorname{id}, \operatorname{id})_{\#}\mu$  is induced by the identity map. In other words,  $\nu = \operatorname{id}_{\#}\mu = \mu$ .

To prove the triangle inequality we will use the gluing lemma below (Lemma 2.3) with N = 3. Let  $\mu_i \in \mathcal{P}_p(X)$  for  $i \in \{1, 2, 3\}$  and let  $\gamma_1 \in \Pi(\mu_1, \mu_2)$  and  $\gamma_2 \in \Pi(\mu_2, \mu_3)$  be optimal in the definition of  $W_p$ . Then, there exists  $\sigma \in \mathcal{P}(X^3)$  such that  $(\pi_{i,i+1})_{\#}\sigma = \gamma_i$  for  $i \in \{1, 2\}$ . A fortiori one has  $(\pi_1)_{\#}\sigma = \mu_1$  and  $(\pi_3)_{\#}\sigma = \mu_3$ , so that  $(\pi_{13})_{\#}\sigma \in \Pi(\mu_1, \mu_3)$ . In particular,

$$\begin{split} W_{p}(\mu_{1},\mu_{3}) &\leqslant \left( \int_{X^{2}} \operatorname{dist}(x,y)^{p} \operatorname{d}(\pi_{1,3})_{\#}\sigma(x,y) \right)^{1/p} \\ &= \left( \int_{X^{3}} \operatorname{dist}(x_{1},x_{3})^{p} \operatorname{d}\sigma(x_{1},x_{2},x_{3}) \right)^{1/p} \\ &\leqslant \left( \int_{X^{3}} (\operatorname{dist}(x_{1},x_{2}) + \operatorname{dist}(x_{2},x_{3}))^{p} \operatorname{d}\sigma(x_{1},x_{2},x_{3}) \right)^{1/p} \\ &\leqslant \left( \int_{X^{3}} \operatorname{dist}(x_{1},x_{2})^{p} \operatorname{d}\sigma(x_{1},x_{2},x_{3}) \right)^{1/p} + \left( \int_{X^{3}} \operatorname{dist}(x_{2},x_{3})^{p} \operatorname{d}\sigma(x_{1},x_{2},x_{3}) \right)^{1/p} \\ &= W_{p}(\mu_{1},\mu_{2}) + W_{p}(\mu_{2},\mu_{3}), \end{split}$$

where we used the Minkowski inequality in  $L^p(\sigma)$  to get the second inequality, and the property  $(\pi_{i,i+1})_{\#}\sigma = \gamma_i$  to get the last equality.

**Lemma 2.3** (Gluing). Let  $X_1, \ldots, X_N$  be complete and separable metric spaces, and for any  $1 \leq i \leq N-1$  consider a transport plan  $\gamma_i \in \Pi(\mu_i, \mu_{i+1})$ . Then, there exists  $\gamma \in \mathcal{P}(X_1, \ldots, X_N)$  such that for all  $i \in \{1, \ldots, N-1\}$ ,  $(\pi_{i,i+1})_{\#}\gamma = \gamma_i$ , where  $\pi_{i,i+1} : X_1 \times \cdots \times X_N \to X_i \times X_{i+1}$  is the projection.

*Proof.* See Lemma 5.3.2 and Remark 5.3.3 in [1].

**Exercise 2.4.** Prove the triangle inequality assuming the existence of optimal transport maps between  $\mu_1, \mu_2$  and  $\mu_2, \mu_3$ .

**Remark 2.5** (Non-compact case). As usual, the compactness assumption is only here for clarity of presentation. In general, when X is a complete and separable metric space, the space  $\mathcal{P}_p(X)$  is defined as the set of probability measures such that for some (and thus any)  $x_0 \in X$  it holds

$$\int \operatorname{dist}(x_0, y)^p \mathrm{d}\mu(y) < \infty.$$

It can be shown that this set endowed with the distance  $W_p$  is also a complete and separable metric space. Exercice: show that the Wasserstein distance  $W_p$  is finite on this set.

#### 2.2 Comparisons

**Comparison between Wasserstein distances** Note that, due to Jensen's inequality, since all  $\gamma \in \Pi(\mu, \nu)$  are probability measures, for  $p \leq q$  we have

$$\left(\int \operatorname{dist}(x,y)^{p} \mathrm{d}\gamma\right)^{\frac{1}{p}} \leqslant \left(\int \operatorname{dist}(x,y)^{q} \mathrm{d}\gamma\right)^{\frac{1}{q}},$$

which implies  $W_p(\mu, \nu) \leq W_q(\mu, \nu)$ . In particular,  $W_1(\mu, \nu) \leq W_p(\mu, \nu)$  for every  $p \geq 1$ . On the other hand, for compact (and thus bounded) X, an opposite inequality also holds, since

$$\left(\int \operatorname{dist}(x,y)^p \mathrm{d}\gamma\right)^{\frac{1}{p}} \leq \operatorname{diam}(X)^{\frac{p-1}{p}} \left(\int \operatorname{dist}(x,y) \mathrm{d}\gamma\right)^{\frac{1}{p}}.$$

This implies that for all  $p \ge 1$ ,

$$W_1(\mu,\nu) \leqslant W_p(\mu,\nu) \leqslant \operatorname{diam}(X)^{\frac{p-1}{p}} W_1(\mu,\nu)^{\frac{1}{p}}.$$

#### 2.3 Topological properties

**Theorem 2.6.** Assume that X is compact. For  $p \in [1, +\infty[$ , we have  $\mu_n \rightarrow \mu$  if and only if  $W_p(\mu_n, \mu) \rightarrow 0$ .

*Proof.* We only need to prove the result for  $W_1$  thanks to the comparison inequalities between  $W_1$  and  $W_p$  in previous section. Let us start from a sequence  $\mu_n$  such that  $W_1(\mu_n,\mu) \to 0$ . Thanks to the duality formula, for every  $\varphi \in \text{Lip}_1(X)$ , we have  $\int \varphi(\mu_n - \mu) \to 0$ . By linearity, the same is true for any Lipschitz function. By density, this holds for any function in  $\mathcal{C}(X)$ . This shows that convergence in  $W_1$  implies weak convergence.

To prove the opposite implication, let us first fix a subsequence  $\mu_{n_k}$  that satisfies  $\lim_k W_1(\mu_{n_k},\mu) = \limsup_n W_1(\mu_n,\mu)$ . For every k, pick a function  $\varphi_{n_k} \in \operatorname{Lip}_1(X)$  such that  $\int \varphi_{n_k}(\mu_{n_k} - \mu) = W_1(\mu_{n_k},\mu)$ . Up to adding a constant, which does not affect the integral, we can assume that the  $\varphi_{n_k}$  all vanish at the same point, and they are hence uniformly bounded and equi-continuous. By Ascoli-Arzelà theorem, we can extract a sub-sequence uniformly converging to a certain  $\varphi \in \operatorname{Lip}_1(X)$ . By replacing the original subsequence with this new one, we have now

$$W_1(\mu_{n_k},\mu) = \int \varphi_{n_k} \mathrm{d}(\mu_{n_k}-\mu) \to \int \varphi \mathrm{d}(\mu-\mu) = 0$$

where the convergence of the integral is justified by the weak convergence  $\mu_{n_k} \rightarrow \mu$  together with the strong convergence in  $\mathcal{C}(X) \ \varphi_{n_k} \rightarrow \varphi$ . This shows that  $\limsup_n W_1(\mu_n, \mu) \leq 0$ and concludes the proof.

**Remark 2.7.** In the non-compact case, it can be shown that convergence in  $\mathcal{P}_p(X)$  is equivalent to tight convergence (in duality with continuous and bounded functions) and convergence of the *p*-th order moments i.e. for all  $x_0 \in X$ ,

$$\int \operatorname{dist}(x_0, y)^p \mathrm{d}\mu_n(y) \to \int \operatorname{dist}(x_0, y)^p \mathrm{d}\mu(y).$$

### 3 Geodesics in Wasserstein space

**Definition 3.1.** Let (X, dist) be a metric space. A constant speed geodesic between two points  $x_0, x_1 \in X$  is a continuous curve  $x : [0,1] \to X$  such that for every  $s, t \in [0,1]$ ,  $\text{dist}(x_s, x_t) = |s - t| \text{dist}(x_0, x_1)$ .

**Proposition 3.2.** Let  $\mu_0, \mu_1 \in \mathcal{P}_p(X)$  with  $X \subset \mathbb{R}^d$  compact and convex. Let  $\gamma \in \Pi(\mu_0, \mu_1)$  be an optimal transport plan. Define

$$\mu_t := (\pi_t)_{\#} \gamma \text{ where } \pi_t(x, y) = (1 - t)x + ty.$$

Then, the curve  $\mu_t$  is a constant speed geodesic between  $\mu_0$  and  $\mu_1$ .

**Example 3.3.** If there exists an optimal transport map T between  $\mu_0$  and  $\mu_1$ , then the geodesic defined above is  $\mu_t = ((1-t)id + tT)_{\#}\mu_0$ .

**Remark 3.4.** In fact, it can be shown that any geodesic between  $\mu_0$  and  $\mu_1$  can be constructed as in Proposition 3.2.

*Proof.* First note that if  $0 \leq s \leq t \leq 1$ ,

$$W_p(\mu_0, \mu_1) \leq W_p(\mu_0, \mu_s) + W_p(\mu_s, \mu_t) + W_p(\mu_t, \mu_1),$$

so that it suffices to prove the inequality  $W_p(\mu_s, \mu_t) \leq |t-s| W_p(\mu_0, \mu_1)$  for all  $0 \leq s \leq t \leq 1$  to get equality. The inequality is easily checked by building an explicit transport plan using an optimal transport plan  $\gamma$ . Take  $\gamma_{st} := (\pi_s, \pi_t)_{\#} \gamma \in \Pi(\mu_s, \mu_t)$ , so that

$$W_{p}(\mu_{s},\mu_{t})^{p} \leq \int ||x-y||^{p} \,\mathrm{d}\gamma_{st}(x,y) = \int ||\pi_{s}(x,y) - \pi_{t}(x,y)||^{p} \,\mathrm{d}\gamma(x,y)$$
$$= \int ||(1-s)x + sy - ((1-t)x + ty)||^{p} \,\mathrm{d}\gamma(x,y)$$
$$= \int ||(t-s)(x-y)||^{p} \,\mathrm{d}\gamma(x,y) = (t-s)^{p} W_{p}(\mu,\nu)^{p} \qquad \Box$$

**Corollary 3.5.** The space  $(\mathcal{P}_p(X), W_p)$  with X compact and convex is a geodesic space, meaning that any  $\mu_0, \mu_1 \in \mathcal{P}_p(X)$  can be joined by (at least one) constant speed geodesic.

## 4 Differentiability of the Wasserstein distance

In this section, we will compute the differential of the Wasserstein distance under additive perturbations.

**Theorem 4.1.** Let  $\sigma$ ,  $\rho_0$ ,  $\rho_1 \in \mathcal{P}(X)$ . Assume that there exists unique Kantorovich potentials  $(\varphi_0, \psi_0)$  between  $\sigma$  and  $\rho_0$  which are c-conjugate to each other and satisfy  $\varphi_0(x_0) = 0$ for some  $x_0 \in X$ . Then,

$$\frac{\mathrm{d}}{\mathrm{d}t} \left| \mathcal{T}_c(\sigma, \rho_0 + t(\rho_1 - \rho_0)) \right|_{t=0} = \int \psi_0 \mathrm{d}(\rho_1 - \rho_0).$$

*Proof.* Denote  $\rho_t = (1-t)\rho_0 + t\rho_1 = \rho_0 + t(\rho_1 - \rho_0)$ . By Kantorovich duality, we have

$$\mathcal{T}_c(\sigma, \rho_t) \ge \int \varphi_0 \mathrm{d}\sigma + \int \psi_0 \mathrm{d}\rho_t.$$

This immediately gives

$$\frac{1}{t}(\mathcal{T}_c(\sigma,\rho_t) - \mathcal{T}_c(\sigma,\rho_0)) \ge \int \psi_0 \mathrm{d}(\rho_1 - \rho_0)$$

To show the converse inequality, we let  $(\varphi_t, \psi_t)$  be *c*-conjugate Kantorovich potentials between  $\sigma$  and  $\rho_t$  satisfying  $\psi_t(x_0) = 0$ , giving

$$\frac{1}{t}(\mathcal{T}_c(\sigma,\rho_0)-\mathcal{T}_c(\sigma,\rho_t)) \ge \int \psi_t \mathrm{d}(\rho_1-\rho_0).$$

Moreover, by uniqueness of  $(\varphi_0, \psi_0)$ , we get that  $\varphi_t, \psi_t$  converges uniformly to  $(\varphi_0, \psi_0)$  as  $t \to 0$ , thus concluding the proof.

The assumption on the uniqueness of the potentials can be guaranteed a priori in the following setting, which corresponds to the distance  $W_2$  (one could prove it for  $W_p$ , with p > 1 similarly).

**Proposition 4.2** (Uniqueness of potentials). If  $X \subseteq \mathbb{R}^d$  is the closure of a bounded and connected open set,  $x_0 \in X$ ,  $(\sigma, \rho) \in \mathcal{P}(X)$  satisfies

$$\operatorname{spt}(\rho) = X \text{ or } \operatorname{spt}(\sigma) = X,$$

then, there exists a unique pair of Kantorovich potentials  $(\varphi, \psi)$  optimal for  $c(x, y) = \frac{1}{2} ||x - y||^2$ , c-conjugate to each other, and satisfying  $\varphi(x_0) = 0$ .

*Proof.* Assume that  $\operatorname{spt}(\sigma) = X$ . Since c is Lipschitz on the bounded set  $X, \varphi, \psi$  are Lipschitz and therefore differentiable almost everywhere. Take  $(x_0, y_0) \in \operatorname{spt}(\gamma)$  where  $\gamma \in \Pi(\sigma, \rho)$  is the optimal transport plan, such that  $\varphi$  is differentiable at  $x_0 \in X$ . As we have already shown, for any optimal pair  $(\varphi, \psi)$  we necessarily have

$$y_0 = x_0 - \nabla \varphi(x_0),$$

so that if  $(\varphi', \psi')$  is another optimal pair, we should have  $\nabla \varphi = \nabla \varphi' \sigma$ -a.e. Since  $\operatorname{spt}(\sigma) = X$  and since X is the closure of a connected open set, this implies  $\varphi = \varphi' + C$  for a constant C as desired, and C = 0 since  $\varphi(x_0) = \varphi'(x_0)$ . Moreover,  $\psi' = \varphi'^c = \varphi^c = \psi$ , allowing to deal with the case where  $\operatorname{spt}(\rho) = X$  by symmetry.

## 5 Dynamic formulation of optimal transport

We conclude this lecture with a discussion around a fluid dynamic interpretation of optimal transport. The material in this section is only treated at an informal level and we refer to [3] for a rigorous treatment.

When  $X \subset \mathbb{R}^d$ , we can interpret the marginals  $\mu, \nu \in \mathcal{P}(X)$  as distributions of particles at times t = 0 and t = 1 respectively. Assume that for each time t, there is a velocity field  $v_t : \mathbb{R}^d \to \mathbb{R}^d$  which moves particles around. The relation between the velocity field and the distribution is given by the continuity equation (satisfied in the sense of distributions)

$$\partial_t \rho_t + \nabla \cdot (\rho_t v_t) = 0.$$

When  $v_t$  is regular enough (e.g. Lipschitz continuous in x, uniformly in t), then we can define its flow  $T : [0,1] \times X \to \mathbb{R}^d$  which is such that  $T_t(x)$  gives the position at time t of a particle which is at x at time 0. It solves  $T_0(x) = x$  and

$$\frac{d}{dt}T_t(x) = v_t(T_t(x)).$$

Let us denote  $CE(\mu, \nu)$  the set of solutions  $(\rho, v)$  to the continuity equation such that  $t \mapsto \rho_t$  is weakly continuous and satisfies  $\rho_0 = \mu$  and  $\rho_1 = \nu$ . Consider also the integrated (generalized) "kinetic energy" functional

$$A_p(\rho, v) := \int_0^1 \int_X \|v_t(x)\|^p \mathrm{d}\mu_t(x) \mathrm{d}t.$$

By minimizing this functional over all interpolations between  $\mu$  and  $\nu$ , we recover the optimal transport with cost  $||y - x||^p$ . This is called the Benamou-Brenier formulation.

**Theorem 5.1** (Dynamic formulation). Let  $\mu, \nu \in \mathcal{P}(\mathbb{R}^d)$  be compactly supported. For  $p \ge 1$  it holds

$$W_p^p(\mu,\nu) = \inf \Big\{ A_p(\rho,v) \mid (\rho,v) \in \operatorname{CE}(\mu,\nu) \Big\}.$$

Let us give some informal arguments to understand this result.

• Let us first argue that for  $(\rho, v) \in CE(\mu, \nu)$  it holds  $A_p(\rho, v) \ge W_p^p(\mu, \nu)$ . Assume  $(\rho, v)$  is regular enough and consider the flow  $T_t(x)$ , that satisfies  $\rho_t = (T_t)_{\#}\rho_0$ . It holds

$$A(\rho, v) = \int_0^1 \int_X \|v_t(T_t(x))\|^p d\rho_0(x) dt$$
  
= 
$$\int_X \left( \int_0^1 \left\| \frac{d}{dt} T_t(x) \right\|^p dt \right) d\rho_0(x)$$
  
$$\geqslant \int_X \|T_1(x) - T_0(x)\|^p d\rho_0(x)$$

by Jensen's inequality. Since  $(T_1)_{\#}\rho_0 = \rho_1 = \nu$  and  $\rho_0 = \mu$ , the last quantity is larger than  $W_p^p(\mu, \nu)$ .

• Let us build an admissible  $(\rho, v) \in CE(\mu, \nu)$  such that  $A(\rho, v) = W_p^p(\mu, \nu)$  using the geodesic between  $\mu$  and  $\nu$ . Assume that there exists an optimal transport map T between  $\mu$  and  $\nu$ , and set  $\rho_t = (T_t)_{\#}\mu$  with  $T_t(x) = (1 - t)x + tT(x)$ . Now define the velocity field

$$v_t = \left(\frac{d}{dt}T_t\right) \circ T_t^{-1} = (T - \mathrm{id}) \circ T_t^{-1},$$

which, by construction, is such that  $(\rho_t, v_t)$  satisfies the continuity equation in the weak sense. We have the desired equality:

$$A(\rho, v) = \int ||v_t(x)||^p \mathrm{d}\rho_t(x) = \int |T(x) - x|^p \mathrm{d}\rho_0(x) = W_p^p(\mu, \nu).$$

**Riemannian interpretation.** In the case p = 2, we can understand (at least at the formal level) the Benamou-Brenier formula as a Riemannian formulation for  $W_2$  (this point of view is due to Otto). In this interpretation, the tangent space at  $\rho \in \mathcal{P}_2(X)$  are measures of the form  $\delta \rho = -\nabla \cdot (v\rho)$  with a velocity field  $v \in L^2(\rho, \mathbb{R}^d)$  and the metric is given by

$$\|\delta\rho\|_{\rho}^{2} = \inf_{v \in L^{2}(\rho, \mathbb{R}^{d})} \left\{ \int \|v(x)\|_{2}^{2} \mathrm{d}\rho(x) \mid \delta\rho = -\nabla \cdot (v\rho) \right\}.$$

## References

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